Observer-based Hamiltonian identification for quantum systems

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Observer-based parameter estimation

Take $\frac{d}{dt}x = f(x, u(t), p)$ with output y(t) = h(x) and control u(t). The goal is to estimate p (and x) from noisy measurements of y.

State-parameter asymptotic observer Can we find g_1 and g_2 such that the solution $(\hat{x}(t), \hat{p}(t))$ of

$$\frac{d}{dt}\hat{x}(t) = f(\hat{x}, u(t), \hat{p}) + g_1(\hat{x}, u(t), \hat{p}, y(t))$$
$$\frac{d}{dt}\hat{p}(t) = g_2(\hat{x}, u(t), \hat{p}, y(t))$$

with an arbitrary initial state (\hat{x}_0, \hat{p}_0) converges towards (x(t), p) as $t \to \infty$?

Invariant asymptotic observers (Thesis of Silvère Bonnabel).

Outline

An invariant asymptotic observer for a 2-level system

Semi-local convergence proof

Possible extensions

The estimation problem for a two-level system

$$\frac{u}{dt}\rho = -\iota \left[\frac{\Delta}{2}\sigma_z + \frac{u\mu}{2}\sigma_x, \rho\right]$$

$$y = \text{Tr}(\sigma_z\rho), \text{ (groud state population)}$$

$$|g\rangle \text{ where:}$$

- ▶ ρ is the density matrix: a 2 × 2 symmetric \geq 0 matrix with $\text{Tr}(\rho) = 1$ and $\text{Tr}(\rho^2) = 1$ (here a projector).
- $\sigma_{x} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_{y} = \begin{pmatrix} 0 & -\iota \\ \iota & 0 \end{pmatrix}, \quad \sigma_{z} = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \text{ are the }$ Pauli matrices: $\sigma_{x}^{2} = 1$; $\sigma_{x}\sigma_{y} = \iota\sigma_{z}$..
- ▶ the two real parameters are \triangle (the difference between the atomic frequency (transition $|g\rangle \leftrightarrow |e\rangle$) and the laser frequency of amplitude u) and $\mu > 0$ the dipole strength.
- \blacktriangleright $u\mu$ is the Rabi pulsation. This model is based on a singular perturbation of a Lindblad equation modelling the evolution of an ensemble of identical 3 level quantum systems.



Hamiltonian identification for quantum systems

Unknown parameters: μ and Δ .

Some previous methods:

Maximum-likelihood (Rabitz, Kosut, Walmsley, Paris, Mabuchi) Maximum (Kullback) entropy (Buzek, Paris, Oliveras) Optimal identification via least-square crieteria (Rabitz, Geremia)

Main issues: Robustness wrt uncertainties and noises, Computational cost, local minima.

Observer-based parameter identification (a previous result):

R.L. Kosut and H. Rabitz, 15th IFAC world congress, 2002. Asymptotic state observer + iterative search algorithm to update the estimate of the parameters

Our approach: adptive observer



The non-linear asymptotic observer

$$\begin{split} \frac{d}{dt}\hat{\rho} &= -\iota \left[\frac{\hat{\Delta}}{2} \sigma_z + \frac{u\hat{\mu}}{2} \sigma_x, \hat{\rho} \right] \\ &- \mathcal{K}_{\rho} (\text{Tr}(\sigma_z \hat{\rho}) - y) \left(\sigma_z \hat{\rho} + \hat{\rho} \sigma_z - 2\text{Tr}(\sigma_z \hat{\rho}) \hat{\rho} \right) \\ \frac{d}{dt}\hat{\mu} &= -u \mathcal{K}_{\mu} \text{Tr}\left(\sigma_y \hat{\rho} \right) \left(\text{Tr}(\sigma_z \hat{\rho}) - y \right) \\ \frac{d}{dt}\hat{\Delta} &= -u \mathcal{K}_{\Delta} \text{Tr}(\sigma_x \hat{\rho}) \left(\text{Tr}(\sigma_z \hat{\rho}) - y \right) \end{split}$$

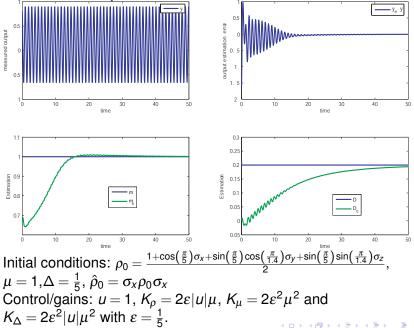
with positive gains K_{ρ} , K_{μ} and K_{Δ} . Preservation of $\text{Tr}(\hat{\rho}) = 1$ and $\text{Tr}(\hat{\rho}^2) = 1$.

Convergence results from averaging consideration (RWA) under the following assumptions and gains design:

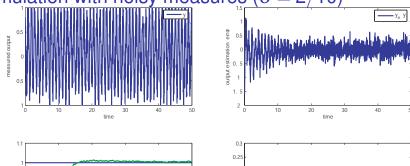
- ▶ slowly varying u versus Rabi pulsation $|u\mu|$: $|\dot{u}| \ll u^2\mu$.
- ▶ Small detuning $|\Delta|, |\hat{\Delta}| \ll |u|\mu$ and $|\hat{\mu}_{t=0} \mu| \ll \mu$.
- ▶ Small gains: $K_{\rho} \ll |u|\mu$, $\sqrt{K_{\mu}} \ll \mu$, $K_{\Delta} \ll K_{\mu}\mu$.

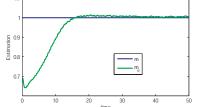


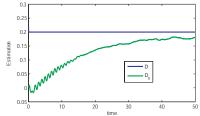
Simulation with perfect measures



Simulation with noisy measures ($\sigma = 2/10$)







Initial conditions:
$$\rho_0 = \frac{1 + \cos(\frac{\pi}{5})\sigma_X + \sin(\frac{\pi}{5})\cos(\frac{\pi}{1.4})\sigma_Y + \sin(\frac{\pi}{5})\sin(\frac{\pi}{1.4})\sigma_Z}{2}$$

 $\mu = 1, \Delta = \frac{1}{5}, \hat{\rho}_0 = \sigma_x \rho_0 \sigma_x$

Control/gains: u=1, $K_{\rho}=2\varepsilon|u|\mu$, $K_{\mu}=2\varepsilon^{2}\mu^{2}$ and

$$K_{\Delta} = 2\varepsilon^2 |u| \mu^2$$
 with $\varepsilon = \frac{1}{5}$.



Invariance versus SU(2) action

For any $U \in SU(2)$, the transformation $((u, y, \Delta, \mu)$ unchanged)

$$\rho \mapsto \overline{\omega} = \mapsto U\rho U^{\dagger}, \sigma_{x} \mapsto \varsigma_{x} = U\sigma_{x} U^{\dagger}, \dots$$

leaves

$$\frac{d}{dt}\rho = -\iota \left[\frac{\Delta}{2}\sigma_z + \frac{u(t)\mu}{2}\sigma_x, \rho\right], \quad y = \operatorname{Tr}(\sigma_z \rho)$$

unchanged:

$$\frac{d}{dt}\varpi = -i\left[\frac{\Delta}{2}\varsigma_{z} + \frac{u(t)\mu}{2}\varsigma_{x},\varpi\right], \quad y = \operatorname{Tr}\left(\varsigma_{z}\varpi\right),$$

and $\zeta_X, \zeta_Y, \zeta_Z$ are new Pauli matrices.

SU(2) invariance of the non-linear observer

For any $U \in SU(2)$, the transformation $((\hat{\Delta}, \hat{\mu}))$ unchanged

$$\hat{\rho} \mapsto \hat{\varpi} = \mapsto U \hat{\rho} U^{\dagger}, \sigma_{x} \mapsto \varsigma_{x} = U \sigma_{x} U^{\dagger}, \dots$$

leaves the asymptotic observer

$$\begin{split} \frac{d}{dt}\hat{\rho} &= -\iota \left[\frac{\hat{\Delta}}{2} \sigma_{z} + \frac{u\hat{\mu}}{2} \sigma_{x}, \hat{\rho} \right] \\ &- \mathcal{K}_{\rho} (\text{Tr}(\sigma_{z}\hat{\rho}) - y) \left(\sigma_{z}\hat{\rho} + \hat{\rho} \sigma_{z} - 2\text{Tr}(\sigma_{x}\hat{\rho})\hat{\rho} \right) \\ \frac{d}{dt}\hat{\mu} &= -u\mathcal{K}_{\mu} \text{Tr}\left(\sigma_{y}\hat{\rho} \right) \left(\text{Tr}(\sigma_{z}\hat{\rho}) - y \right) \\ \frac{d}{dt}\hat{\Delta} &= -u\mathcal{K}_{\Delta} \text{Tr}(\sigma_{x}\hat{\rho}) \left(\text{Tr}(\sigma_{z}\hat{\rho}) - y \right) \end{split}$$

unchanged.



Assumptions

In

$$\begin{split} \frac{d}{dt}\hat{\rho} &= -\iota \left[\frac{\hat{\Delta}}{2} \sigma_{z} + \frac{u\hat{\mu}}{2} \sigma_{x}, \hat{\rho} \right] \\ &- \mathcal{K}_{\rho} (\text{Tr}(\sigma_{z}\hat{\rho}) - y) \left(\sigma_{z}\hat{\rho} + \hat{\rho} \sigma_{z} - 2\text{Tr}(\sigma_{z}\hat{\rho})\hat{\rho} \right) \\ \frac{d}{dt}\hat{\mu} &= -u\mathcal{K}_{\mu} \text{Tr}\left(\sigma_{y}\hat{\rho} \right) \left(\text{Tr}(\sigma_{z}\hat{\rho}) - y \right) \\ \frac{d}{dt}\hat{\Delta} &= -u\mathcal{K}_{\Delta} \text{Tr}(\sigma_{x}\hat{\rho}) \left(\text{Tr}(\sigma_{z}\hat{\rho}) - y \right) \end{split}$$

we assume that u is constant and that

$$\hat{\Delta} = \varepsilon \hat{d}, \quad K_{\rho} = 4k_{\rho} \varepsilon |u| \mu, \quad K_{\mu} = 2k_{\mu} \varepsilon^{2} \mu^{2}, \quad K_{\Delta} = 2k_{\Delta} \varepsilon^{2} |u| \mu^{2}$$

for $\varepsilon>0$ small $\varepsilon\ll 1,\, k_{\rho},k_{\mu},k_{\Delta}\sim 1.$

Convergence based on perturbation techniques (Rotating Wave Approximation (RWA)) but up to order 2 in ε .



In the interaction frame

Consider the following time-varying transformation

$$\rho = e^{-\imath \frac{\iota \mu \iota \tau \sigma_X}{2}} \xi \, e^{\imath \frac{\iota \mu \iota \sigma_X}{2}}, \qquad \hat{\rho} = e^{-\imath \frac{\iota \mu \iota \tau \sigma_X}{2}} \hat{\xi} \, e^{\imath \frac{\iota \mu \iota \tau \sigma_X}{2}}.$$

The dynamics reads:

$$\begin{split} \frac{d}{dt} & \hat{\xi} = - \imath \left[\frac{\Delta}{2} e^{\imath \iota \mu t \sigma_x} \sigma_z, \xi \right] \\ \frac{d}{dt} & \hat{\xi} = - \imath \left[\frac{\hat{\Delta}}{2} e^{\imath \iota \mu t \sigma_x} \sigma_z + \frac{\iota \iota (\hat{\mu} - \mu)}{2} \sigma_x, \hat{\xi} \right] - \mathcal{K}_\rho \text{Tr} \left(e^{\imath \iota \mu t \sigma_x} \sigma_z (\hat{\xi} - \xi) \right) \\ & \qquad \qquad \dots \left(e^{\imath \iota \iota \mu t \sigma_x} \sigma_z \hat{\xi} + \hat{\xi} e^{\imath \iota \iota \mu t \sigma_x} \sigma_z - 2 \text{Tr} \left(e^{\imath \iota \iota \mu t \sigma_x} \sigma_z \hat{\xi} \right) \hat{\xi} \right) \\ \frac{d}{dt} & \hat{\mu} = - \iota \mathcal{K}_\mu \text{Tr} \left(e^{\imath \iota \iota \mu t \sigma_x} \sigma_y \hat{\xi} \right) \text{Tr} \left(e^{\imath \iota \iota \mu t \sigma_x} \sigma_z (\hat{\xi} - \xi) \right) \\ \frac{d}{dt} & \hat{\Delta} = - \iota \mathcal{K}_\Delta \text{Tr} \left(\sigma_x \hat{\xi} \right) \text{Tr} \left(e^{\imath \iota \iota \mu t \sigma_x} \sigma_z (\hat{\xi} - \xi) \right). \end{split}$$

First order secular approximation

By assumption the frequency $u\mu$ is large and the integration of $e^{iu\mu t\sigma_x}$ will produce terms of small amplitude. We neglect terms rotating at $u\mu$ and also $2u\mu$ (first order in ε):

$$\begin{split} \frac{d}{dt} \xi &= 0 \\ \frac{d}{dt} \hat{\xi} &= -\imath \, \left[\frac{u(\hat{\mu} - \mu)}{2} \sigma_x, \hat{\xi} \right] \\ &\quad - \frac{K_\rho}{2} \text{Tr} \left(\sigma_y (\hat{\xi} - \xi) \right) \left(\sigma_y \hat{\xi} + \hat{\xi} \sigma_y - 2 \text{Tr} \left(\sigma_y \hat{\xi} \right) \hat{\xi} \right) \\ &\quad - \frac{K_\rho}{2} \text{Tr} \left(\sigma_z (\hat{\xi} - \xi) \right) \left(\sigma_z \hat{\xi} + \hat{\xi} \sigma_z - 2 \text{Tr} \left(\sigma_z \hat{\xi} \right) \hat{\xi} \right) \\ &\quad \frac{d}{dt} \hat{\mu} = - \frac{u K_\mu}{2} \left(\text{Tr} \left(\sigma_y \hat{\xi} \right) \text{Tr} \left(\sigma_z (\hat{\xi} - \xi) \right) - \text{Tr} \left(\sigma_z \hat{\xi} \right) \text{Tr} \left(\sigma_y (\hat{\xi} - \xi) \right) \right) \\ \frac{d}{dt} \hat{\Delta} &= 0. \end{split}$$

Convergence of $\hat{\xi}$ and $\hat{\mu}$

Up to second order terms, $\hat{\xi}$ and $\hat{\mu}$ obey an autonomous differential system where ξ is a parameter:

$$\begin{split} \frac{d}{dt}\hat{\xi} &= -\imath \, \left[\frac{u(\hat{\mu} - \mu)}{2} \sigma_x, \hat{\xi} \right] \\ &- \frac{K_\rho}{2} \text{Tr} \left(\sigma_y(\hat{\xi} - \xi) \right) \left(\sigma_y \hat{\xi} + \hat{\xi} \sigma_y - 2 \text{Tr} \left(\sigma_y \hat{\xi} \right) \hat{\xi} \right) \\ &- \frac{K_\rho}{2} \text{Tr} \left(\sigma_z(\hat{\xi} - \xi) \right) \left(\sigma_z \hat{\xi} + \hat{\xi} \sigma_z - 2 \text{Tr} \left(\sigma_z \hat{\xi} \right) \hat{\xi} \right) \\ \frac{d}{dt} \hat{\mu} &= - \frac{u K_\mu}{2} \left(\text{Tr} \left(\sigma_y \hat{\xi} \right) \text{Tr} \left(\sigma_z (\hat{\xi} - \xi) \right) - \text{Tr} \left(\sigma_z \hat{\xi} \right) \text{Tr} \left(\sigma_y (\hat{\xi} - \xi) \right) \right) \end{split}$$

Local exponential convergence for any ξ (excepted some isolated values) and for any K_{ρ} , $K_{\mu}>0$ via the Lyapounov function:

$$\frac{1}{2}\text{Tr}\left(\sigma_{\!\scriptscriptstyle \mathcal{Y}}(\hat{\xi}-\xi)\right)^2 + \frac{1}{2}\text{Tr}\left(\sigma_{\!\scriptscriptstyle \mathcal{Z}}(\hat{\xi}-\xi)\right)^2 + \frac{1}{K_{\!\scriptscriptstyle \mathcal{U}}}(\hat{\mu}-\mu)^2.$$

Second order secular approximation

We use Kapitsa short-cut method to compute these second order terms particularly important for ξ and $\hat{\Delta}$ since the first order secular terms vanish.

We can decompose $\xi = \bar{\xi} + \delta \xi$: $\bar{\xi}$ is the no-oscillatory part, whereas $\delta \xi$ is the oscillatory one with zero mean and small amplitude $\|\delta\xi\| \ll \|\bar{\xi}\|$. Since $\frac{d}{dt}\xi = -i \left[\frac{\Delta}{2}e^{iu\mu t\sigma_x}\sigma_z, \xi\right]$ we have approximatively:

$$\delta \xi = \frac{\imath \Delta}{2 u \mu} \left[\frac{\Delta}{2} e^{\imath u \mu t \sigma_x} \sigma_y, \bar{\xi} \right] + \dots$$

Plugging this relation into the true dynamics of ξ and taking the secular terms yields up to the third order:

$$\frac{d}{dt}\xi = -i\frac{\Delta^2}{2u\mu}[\sigma_x, \xi] + \dots$$

the term $\frac{\Delta^2}{2\mu\mu}$ corresponds exactly to Bloch-Siegert frequency shift.



Second order secular approximation (continued)

Since $\frac{d}{dt}\hat{\Delta} = -uK_{\Delta}\mathrm{Tr}\left(\sigma_{x}\hat{\xi}\right)\,\mathrm{Tr}\left(e^{\iota u\mu t\sigma_{x}}\sigma_{z}(\hat{\xi}-\xi)\right)$ the secular effect can only comes from the part of $\delta\xi$ and $\delta\hat{\xi}$ with frequency $u\mu$: terms of double frequency $2u\mu$ have no secular effect. In the $\hat{\xi}$ dynamics

$$\begin{split} \frac{d}{dt}\hat{\xi} &= -i \left[\frac{\hat{\Delta}}{2} e^{iu\mu t \sigma_{x}} \sigma_{z} + \frac{u(\hat{\mu} - \mu)}{2} \sigma_{x}, \hat{\xi} \right] - K_{\rho} \text{Tr} \left(e^{iu\mu t \sigma_{x}} \sigma_{z} (\hat{\xi} - \xi) \right) \\ & \dots \left(e^{iu\mu t \sigma_{x}} \sigma_{z} \hat{\xi} + \hat{\xi} e^{iu\mu t \sigma_{x}} \sigma_{z} - 2 \text{Tr} \left(e^{iu\mu t \sigma_{x}} \sigma_{z} \hat{\xi} \right) \hat{\xi} \right) \end{split}$$

the $u\mu$ frequency oscillatory term, denoted by $\delta_1\hat{\xi}$, comes only from $-\iota$ $\left[\frac{\hat{\Delta}}{2}e^{\iota u\mu t\sigma_x}\sigma_z,\hat{\xi}\right]$. Thus

$$\delta_1 \hat{\xi} = \frac{\imath \hat{\Delta}}{2\mu\mu} \left[e^{\imath u\mu t \sigma_x} \sigma_y, \hat{\xi} \right] \quad \text{and} \quad \delta \xi = \delta_1 \xi = \frac{\imath \Delta}{2\mu\mu} \left[e^{\imath u\mu t \sigma_x} \sigma_y, \xi \right]$$

Second order secular approximation (end)

We have the following triangular and locally convergent dynamics:

$$\begin{split} \frac{d}{dt}\hat{\xi} \stackrel{\text{order 1}}{=} - \iota \left[\frac{u(\hat{\mu} - \mu)}{2} \sigma_{x}, \hat{\xi} \right] \\ - \frac{K_{\rho}}{2} \text{Tr} \left(\sigma_{y}(\hat{\xi} - \xi) \right) \left(\sigma_{y}\hat{\xi} + \hat{\xi} \sigma_{y} - 2 \text{Tr} \left(\sigma_{y}\hat{\xi} \right) \hat{\xi} \right) \\ - \frac{K_{\rho}}{2} \text{Tr} \left(\sigma_{z}(\hat{\xi} - \xi) \right) \left(\sigma_{z}\hat{\xi} + \hat{\xi} \sigma_{z} - 2 \text{Tr} \left(\sigma_{z}\hat{\xi} \right) \hat{\xi} \right) \\ \frac{d}{dt}\hat{\mu} \stackrel{\text{order 1}}{=} - \frac{uK_{\mu}}{2} \left(\text{Tr} \left(\sigma_{y}\hat{\xi} \right) \text{Tr} \left(\sigma_{z}(\hat{\xi} - \xi) \right) - \text{Tr} \left(\sigma_{z}\hat{\xi} \right) \text{Tr} \left(\sigma_{y}(\hat{\xi} - \xi) \right) \right) \\ \frac{d}{dt}\xi \stackrel{\text{order 2}}{=} - \iota \frac{\Delta^{2}}{2uu} [\sigma_{x}, \xi] \end{split}$$

$$\begin{split} \frac{d}{dt} \hat{\Delta} & \stackrel{\text{order 2}}{=} -\frac{K_{\Delta}}{\mu} \left(\text{Tr} \left(\sigma_{x} \hat{\xi} \right)^{2} \hat{\Delta} - \text{Tr} \left(\sigma_{x} \hat{\xi} \right) \text{Tr} (\sigma_{x} \xi) \Delta \right) \\ & + \frac{K_{\Delta} \hat{\Delta}}{2\mu} \left(\text{Tr} \left(\sigma_{y} \hat{\xi} \right) \text{Tr} \left(\sigma_{y} (\hat{\xi} - \xi) \right) - \text{Tr} \left(\sigma_{z} \hat{\xi} \right) \text{Tr} \left(\sigma_{z} (\hat{\xi} - \xi) \right) \right). \end{split}$$

Gain design via linear tangent approximation With

$$\hat{\xi} - \xi = \frac{1 + \tilde{x}\sigma_x + \tilde{y}\sigma_y + \tilde{z}\sigma_z}{2}, \quad \tilde{\mu} = \hat{\mu} - \mu, \quad \tilde{\Delta} = \hat{\Delta} - \Delta$$

we have, around $\rho = \frac{1-\sigma_z}{2}$;

$$\frac{d}{dt}\tilde{y} = -u\tilde{\mu} - K_{\rho}\tilde{y}, \qquad \frac{d}{dt}\tilde{\mu} = uK_{\mu}\tilde{y}/2$$

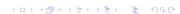
and around $\rho = \frac{1-\sigma_{\chi}}{2}$

$$\frac{d}{dt}\tilde{\Delta} = -\frac{K_{\Delta}}{\mu}\tilde{\Delta}.$$

To respect the scaling, choose $0 < \varepsilon \ll 1$ and set

$$K_{\rho} = 2k_{\rho}\varepsilon|u|\mu, \quad K_{\mu} = 2\varepsilon^{2}\mu^{2}, \quad K_{\Delta} = k_{\Delta}\varepsilon^{2}|u|\mu^{2}$$

with k_o, k_Δ around 1.



N-level system

The system is ($\Delta^{kl} = 0$, no laser de-tuning here)

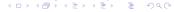
$$\frac{d}{dt}\rho = -\iota \left[\sum_{kl} \frac{u^{kl}\mu^{kl}}{2} \sigma_x^{kl}, \rho\right], \quad y_k = \operatorname{Tr}(P_k\rho)$$

with $P_k = |k\rangle \langle k|$ and its asymptotic observer reads:

$$\begin{split} \frac{d}{dt}\hat{\rho} &= -\iota \left[\sum_{kl} \frac{u^{kl} \hat{\mu}^{kl}}{2} \sigma_x^{kl}, \hat{\rho} \right] \\ &- \sum_k K_\rho^k (\text{Tr}(P_k \hat{\rho}) - y_k) \, \left(P_k \hat{\rho} + \hat{\rho} P_k - 2 \text{Tr}(P_k \hat{\rho}) \hat{\rho} \right) \\ \frac{d}{dt} \hat{\mu}^{kl} &= - \sum_{kl} K_\mu^{kl} \text{Tr} \left(u \sigma_y^{kl} \hat{\rho} \right) \, \left(\text{Tr} \left(\sigma_z^{kl} \hat{\rho} \right) - y_k + y_l \right) \end{split}$$

where $\sigma_{x}^{kl} = |k\rangle \langle l| + |l\rangle \langle k|, \ldots$

Such extensions are possible since we start with an invariant observer for the 2-level system, i.e. we exploit the geometry.



Previous works

- Identifiability for quantum systems: see, e.g., C. Lebris et al (COCV) where it is shown that resonant controls are sufficient.
- ► Asymptotic observers and symmetries: few references (Aghannan, Bonnabel, Martin, R., Dayawansa and coworkers). See the preprint on Symmetry preserving Observers: http://arxiv.org/abs/math.0C/0612193 IEEE TAC Jan 2009

Measurement process